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A new PET prototype for proton therapy: comparison of data and Monte Carlo simulations

V. Rosso,^{a,1} G. Battistoni,^b N. Belcari,^a N. Camarlinghi,^a A. Ferrari,^c S. Ferretti,^a
A. Kraan,^a A. Mairani,^d N. Marino,^a J. E. Ortuño,^{e,f} M. Pullia,^d P. Sala,^b A. Santos,^{e,f}
G. Sportelli,^a K. Straub^a and A. Del Guerra^a

^a*Department of Physics, University of Pisa, and INFN Sezione di Pisa,
Largo Bruno Pontecorvo, 3, 56127 Pisa, Italy*

^b*INFN Sezione di Milano,
Via Celoria 16, 20133 Milano, Italy*

^c*CERN,
1211 Geneve 23, Switzerland*

^d*Fondazione CNAO,
Pavia, Italy*

^e*Biomedical Image Technologies, E.T.S.I.T., Universidad Politécnica de Madrid,
Ciudad Universitaria,, 28040 Madrid, Spain*

^f*CIBER de Bioingeniería, Biomateriales y Nanomedicina (CIBER-BBN),
Ciudad Universitaria, 28040 Madrid, Spain*

E-mail: valeria.rosso@pi.infn.it

¹Corresponding author.

ABSTRACT: Ion beam therapy is a valuable method for the treatment of deep-seated and radio-resistant tumors thanks to the favorable depth-dose distribution characterized by the Bragg peak. Hadrontherapy facilities take advantage of the specific ion range, resulting in a highly conformal dose in the target volume, while the dose in critical organs is reduced as compared to photon therapy. The necessity to monitor the delivery precision, i.e. the ion range, is unquestionable, thus different approaches have been investigated, such as the detection of prompt photons or annihilation photons of positron emitter nuclei created during the therapeutic treatment. Based on the measurement of the induced β^+ activity, our group has developed various in-beam PET prototypes: the one under test is composed by two planar detector heads, each one consisting of four modules with a total active area of $10 \times 10 \text{ cm}^2$. A single detector module is made of a LYSO crystal matrix coupled to a position sensitive photomultiplier and is read-out by dedicated frontend electronics. A preliminary data taking was performed at the Italian National Centre for Oncological Hadron Therapy (CNAO, Pavia), using proton beams in the energy range of 93–112 MeV impinging on a plastic phantom. The measured activity profiles are presented and compared with the simulated ones based on the Monte Carlo FLUKA package.

KEYWORDS: Instrumentation for hadron therapy; Instrumentation for heavy-ion therapy; Gamma camera, SPECT, PET PET/CT, coronary CT angiography (CTA)

Contents

1	Introduction	1
2	Methods and materials	2
2.1	Experimental setup	2
2.2	Reconstruction algorithm	3
2.3	Data analysis and simulations	3
3	Results and discussion	3
4	Conclusion	5

1 Introduction

Proton therapy is a treatment modality of great interest in radiation oncology, because it allows for highly conformal dose distribution and improved sparing of critical structures as compared to treatments with photons. However, proton treatments are much more sensitive to uncertainties such as proton range errors, patient positioning errors, and anatomical variations compared to radiotherapy with photons. As a result, the dose delivered with protons may be significantly different from what was planned, hereby possibly erasing the beneficial effects of proton therapy. To ensure a successful dose delivery, it is highly desirable to have a direct verification of the particle range and delivered dose.

One of the most promising techniques for this purpose is Positron Emission Tomography (PET). PET imaging is based on the detection of the two annihilation gamma-rays following the decay of β^+ isotopes created during irradiation with hadrons. Non-invasive dose monitoring is achieved by comparing the measured and pre-calculated β^+ activity distributions, during and/or after therapeutic irradiation. PET is already successfully used in cancer treatments with particles [1, 2] to verify the ion range. However the distal fall-off of the activity distribution from proton irradiation cannot directly be matched to the dose distribution, making it a challenging task. Information about the delivered dose can be obtained by measuring the activity profile, which are indirectly correlated [3]. In [4], we presented a PET prototype developed at the INFN and the Physics Department of the University of Pisa with an active area of $5 \times 5 \text{ cm}^2$, which already proved its capability to acquire data under clinical conditions. Here, we test the performance of a new, larger, PET prototype with an active area of $10 \times 10 \text{ cm}^2$ regarding the determination of the activity distal fall-off position in a PMMA target, irradiated with mono-energetic protons of energies in the range 93–112 MeV. This is done by comparing the reconstructed β^+ activity distributions from data taken at the Centro Nazionale di Adroterapia Oncologica CNAO [5] (Italian National Center for Oncological Hadron Therapy) with Monte Carlo predictions from the FLUKA package [6, 7].

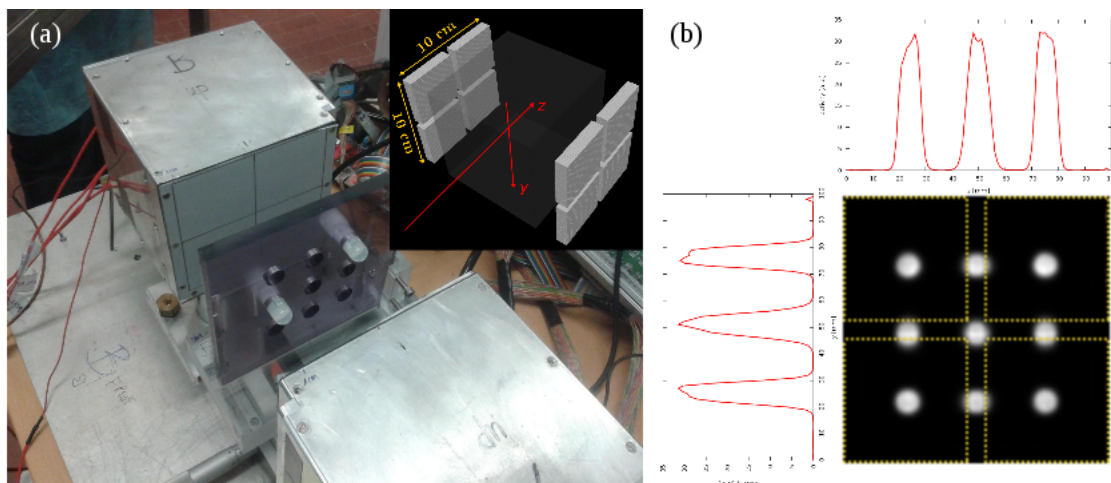


Figure 1. (a) The PET detector system and the phantom with two out of 9 FDG-filled test tubes mounted. Each tube has an outer diameter of 7 mm; the distance between two test tubes is 26 mm. The layout of the phantom was chosen so that the tubes are positioned either in the centre or along the gap between the four detector modules. The inset shows a schematic view of the PET detector system. Each of the two stationary heads contains four independent, closely mounted detector modules with an active area of roughly $5 \times 5 \text{ cm}^2$ each. The 6 mm gaps between the modules are visible. The beam direction is parallel to the z axis. (b) Reconstructed image averaging 20 central slices in the $y - z$ plane and the corresponding profiles along the y and z direction. The software allows reconstructing the activity contained in the tubes independently from their positioning along the $y - z$ plane.

2 Methods and materials

2.1 Experimental setup

The compact PET detector system developed at the INFN and the University of Pisa is designed to be compatible with the installation on a gantry for beam delivery for in-beam data taking. It is composed of two stationary planar heads with an active area of $10 \times 10 \text{ cm}^2$, positioned 14 cm apart, see figure 1. The best spatial resolution is achieved along the beam (z) axis to optimize the determination of the distal fall-off of the induced activity. The plastic phantoms were centred in between the two heads, irradiated by a horizontally impinging proton beam. The new device followed closely the layout of our former PET prototype ($5 \times 5 \text{ cm}^2$), already described in detail in [4]. In short, each head contains four detector modules, made of a position sensitive photomultiplier coupled to a scintillating LYSO crystal matrix (23×23 pixels, pitch 2.0 mm). Each module is read out by dedicated front-end electronics, based on short deadtime pulse discriminators and a series of peak detectors providing position and energy signals. A modular acquisition system has been used for detecting and storing the coincidences between opposing heads [9] with a coincidence window of 10 ns. The energy resolution is 18% at 511 keV (FWHM). The acquisition coincidence rate was in the range of a few kHz.

Data were taken at the CNAO therapy centre. A detailed explanation of the CNAO proton beamline is given in [8]. We used monoenergetic proton pencil beams of 93, 97.5, 99, and 112 MeV with a beam diameter (FWHM) of 14 mm. A PMMA target ($\text{C}_5\text{H}_8\text{O}_2$, density: 1.18 g/cm^3) of $5 \times$

$5 \times 10 \text{ cm}^3$ was irradiated with 5×10^9 protons for about 30 seconds. The experimental data were acquired for 20 minutes starting directly after the irradiation of the target (off-beam measurement).

2.2 Reconstruction algorithm

To reconstruct the acquired coincidences a fully three-dimensional list-mode Maximum Likelihood Expectation Maximization (MLEM) algorithm was implemented. For a detailed description of the algorithm the reader is referred to [10].

The system response probabilities needed for the MLEM are calculated on-the-fly using elliptic tubes of response whose parameters are derived from Monte Carlo simulation. The software makes use of massive parallelism, therefore only few seconds per iteration are needed to reconstruct one acquisition. A total of five iterations are carried out for each acquisition as no further improvements were visually detected after the 5th iteration. The reconstructed images cover a Field of View (FOV) of $10 \times 10 \times 10 \text{ cm}^3$ with an isotropic voxel size of 1 mm.

2.3 Data analysis and simulations

The data analysis focused on the determination of the distal fall-off of the β^+ activity profiles acquired for 20 minutes after the end of proton irradiation.

Starting from both experimental and simulated activity distributions, the activity profiles were made by summing all events within a cylinder with its centre corresponding to the beam axis and projecting the sum of each slice onto the z (beam) axis. To determine the position μ of the distal fall-off at the 50% level of the activity profile along z , a complementary error function was fitted to the data (see section 3, figure 2(b)). After having determined the proximal rising edge at the 50% level, it was used to set the $z = 0$ mm.

Monte Carlo simulations were performed using the FLUKA development version, providing information on the complete transport of the primary protons and the produced secondary particles. We have applied the default settings for hadron therapy. The geometry used in the simulation code consisted of the CNAO proton beamline [4] and a thick PMMA target of $5 \times 5 \times 10 \text{ cm}^3$. We scored the β^+ activity through the USRBIN parameter ACTIVITY at a time $t = 20$ minutes after proton irradiation. According to the simulation the main activity contribution within 20 minutes after proton irradiation comes from the β^+ decaying isotopes ^{15}O ($T_{1/2} = 122$ sec) and ^{11}C ($T_{1/2} = 20.39$ min).

3 Results and discussion

Due to its design made out of four independent detector modules in each head, approximately 6 mm narrow non-active gaps between the modules show up. To test the imaging capabilities of the PET detector system, a phantom based on several test tubes was realized, see figure 1(a). Each tube is positioned so as to be either in the centre of one or along the gap between the four detector modules. A total of nine test tubes were filled with the same FDG activity and put into a PMMA support. The final reconstructed image averaging 20 central slices in the $y - z$ plane is shown in figure 1(b) as well as the profiles taken along the y and z directions. This quality test shows that the software allows reconstructing the activity contained in the tubes independently from their positioning along the $y - z$ plane.

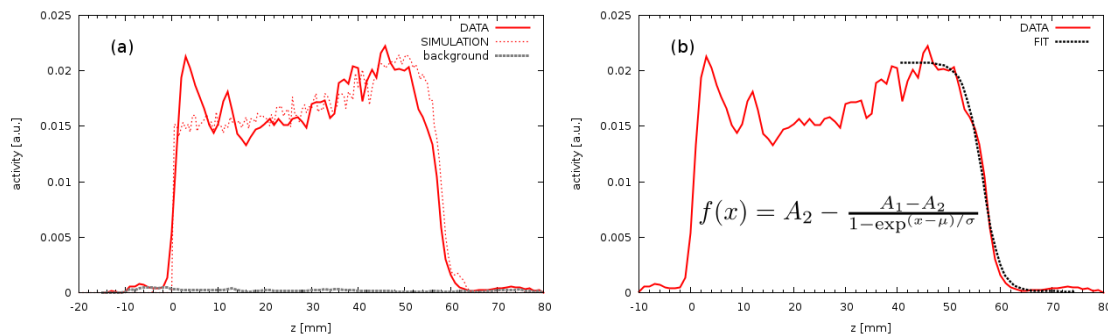


Figure 2. Figure (a) shows the normalized experimental and simulated activity profile along the z (beam) axis acquired for 20 minutes after proton irradiation for a proton energy of 99 MeV. The positions of the 50% distal fall-off agree quite well. The dotted line corresponds to a constant reconstructed background can be seen all over the reconstructed volume. In figure (b) the complementary error function $f(x)$ was fitted to the 99 MeV data yielding the plateau and base value of the profile (A_1 and A_2 , respectively), the position of the 50% distal fall-off (μ) and the slope (σ) of the distribution.

As a result of the analysis on the data, taken at the CNAO therapy centre, figure 2(a) shows the non-quantitative activity profile along the beam (z) axis for the 99 MeV beam energy data set as well as the corresponding simulated profile 20 minutes after proton irradiation of a PMMA target.

The noise which is seen for $z < 0$ mm and after the distal fall-off is compatible with the reconstructed activity background which is shown in figure 2(a), dotted line. The activity background profile was taken from a region outside the irradiated phantom volume and is constant all over the reconstructed FOV.

For a direct comparison, the background corrected area under the profiles is normalized to 1. Despite small oscillations in the plateau region due to a preliminary reconstruction model for the experimental data, the figure shows that the positions of the fall-off for data and simulation agree within 1 mm, which was also verified for the other proton energies.

In figure 2(b) we show an example of the 99 MeV data fitted with the complementary error function returning the mean value μ and standard deviation of the distal 50% *fall-off* depths. The mean value μ was tested to be stable within 1 mm by varying the fitting region in z between 0 and 80 mm.

Results reported from now on are obtained using a fitting region of 30 mm around the activity distal fall-off. The dependence of the fit on the counting statistics was studied by analysing subsets of the entire data acquisition. This was done by varying the diameter of the cylinder (see subsection 2.2) from 5 to 25 mm in 5 mm steps. For each profile, the fit yielded a robust value for μ , changing less than 1 mm. A diameter of 15 mm was chosen as default value, close to the beam diameter.

Figure 3 summarizes the accuracy of our PET prototype concerning the position of the activity profile fall-off for monoenergetic protons in PMMA. A good agreement between data and simulations is seen, showing the capability of our new PET prototype to determine the proton fall-off position with a precision of roughly 1 mm. This is about half the pixel size of the scintillator matrix, which is 2.0 mm.

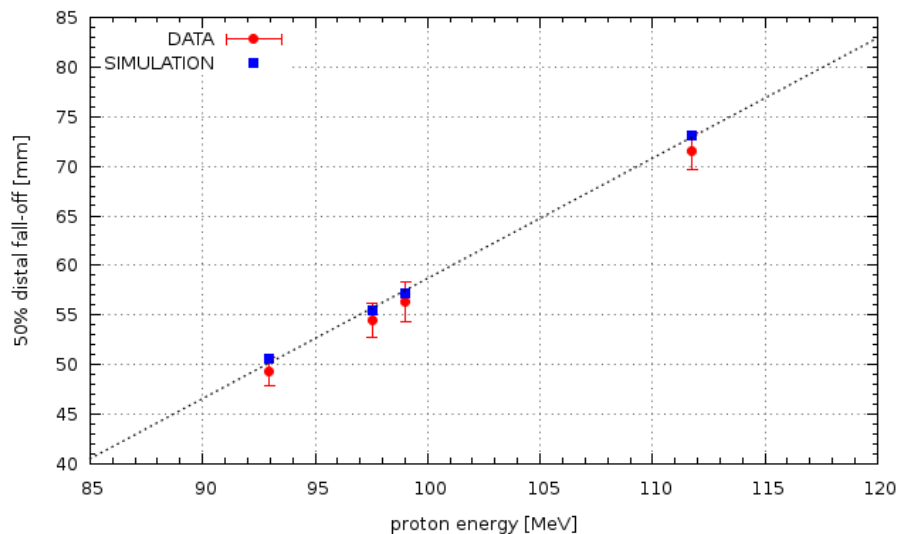


Figure 3. Comparison of the position of the 50% fall-off between FLUKA simulation and data for the energies 93, 97.5, 99 and 112 MeV. We determine a 1 mm accuracy for the determination of the distal fall-off for our $10 \times 10 \text{ cm}^2$ PET prototype. The error bar represents the σ value of the fit. (We assume a systematic error of 1 mm for all energies coming from the positioning of the phantom with respect to the detector heads, which is not shown here.)

4 Conclusion

In this study we have for the first time tested the performance of a new $10 \times 10 \text{ cm}^2$ PET prototype in terms of the activity profile fall-off position for monoenergetic protons between 93 and 112 MeV impinging on a homogeneous PMMA target. The precision of the system to determine the proton fall-off position is of the order of 1 mm. The reconstruction algorithm is still under development and will be refined. The achieved spatial resolution is sufficient to extend this study to both multi-material and antropomorphic phantoms. As already demonstrated with our previous $5 \times 10 \text{ cm}^2$ PET prototype, full in-beam measurements are feasible [11]. An analysis of the in-beam data for the $10 \times 10 \text{ cm}^2$ prototype is currently being performed and will be presented in a future work.

Acknowledgments

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